

HYPOTHETICAL APPROACH IN DETERMINING VIBRATIONS OF PERIODIC CUTTING TOOL HOLDER

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ABSTRACT

Vibration is a standout amongst the most irritating issues looked amid the metal cutting activity, and it happens much of the time in assembling ventures. The vibration level relies upon a wide range of parameters, for example, material sort, inflexibility of tooling structure, cutting information and task mode. In processing, the slicing procedure exposed to the device vibrations having a processing tool holder will doubtlessly result in high vibration levels. These vibrations have an outcome of diminished tool life, poor surface complete and sound disseminations. This examination shows another methodology of confinement for a versatile occasional cutting tool holder of processing machine. A numerical model has been produced to portray the structure of the cutting tool holder. Then again, the conduct of occasional holder is explored numerically. This paper inspected the overwhelming processing vibration parts and recognized these vibrations, which are identified with auxiliary powerful properties of the processing periodic tool holder.

KEYWORDS: Milling; Vibration; Modelling, Periodic Holder

INTRODUCTION

The fundamental thought basic the entire idea of occasional structures is that when a wave is going in a medium and meets a progress in those medium attributes, apiece of it will engender through the new medium and another part will reflect. While, in a customary structure, the wave is relied upon to go with no change until the point that it achieves the limits of the structure. The capacity of periodic structures to transmit waves starting with one area then onto the next inside the pass groups can be significantly decreased when the perfect periodicity is disturbed or scattered.

A wave propagation based approach for the detection of damage in components of structures having periodic damage has been proposed. Periodic damage pattern may arise in a structure due to periodicity in geometry and in loading, Mukherjee [1]. Christoph Ertelt, [2] presented an approach to unify knowledge for generative design and generative fabrication. The geometry of designs and their mapping to removal volumes corresponding to fabrication processes on CNC machine tools are represented. Maria, [3] presented a mathematical model for the propagation of structural waves on an infinitely long, periodically supported beam. The wave types that can exist on the beam are bending waves with displacements in the horizontal and vertical directions, compression waves and torsional waves. A piece of the reflected wave will cooperate with the occurrence wave in a way that will portray the obstruction. At the point when productive obstruction happens, the recurrence is described by being the pass band of the structure, while, because of ruinous impedance, the recurrence is portrayed by being the stop band of the structure. Mechanical instabilities in periodic porous

elastic structures may lead to the formation of homogeneous patterns, opening avenues for a wide range of applications that are related to the geometry of the system, Sicong Shan [4].

Shin, [5] introduced a frequency-domain method of structural damage identification. It is formulated in a general form from the dynamic stiffness equation of motion for a structure and then applied to a beam structure. Only the dynamic stiffness matrix for the intact state appears in the final form of the damage identification algorithm as the structure model. On the off chance, that the structure setup is rehashed for a few times, it is known as a periodic structure. The ruinous impacts will indicate more fundamentally, when the redundancies of the structure unit increment in number, in light of the fact that as the piece of the wave that proliferates consolidates other comparative changes in the medium, another piece of it is destructed, etc.

Many structural components can be regarded as waveguides. They are uniform in one direction so that the cross section of the waveguide has the same physical and geometric properties at all points along the axis of the waveguide. Duhamel [6] presented a method to calculate the forced response of such a structure using a combination of wave and finite element (FE) approaches. The method involves post-processing a conventional, but low order, FE model in which the mass and stiffness matrices are typically found using a conventional FE package. Macea [7] described a method by which the dispersion relations for a two-dimensional structural component can be predicted from a finite element (FE) model. The structure is homogeneous in two dimensions but the properties might vary through the thickness. Toward the application of a two-scale analysis method for nonlinear heterogeneous solids with periodic microstructures, we make a study and introduce a parallel algorithm to achieve the computational efficiency, Matsui [8]. Instances of occasional structures can be found in oil pipelines, railroad tracks, and numerous others. An outline of a straightforward periodic device holder framework is appeared in figure 1.

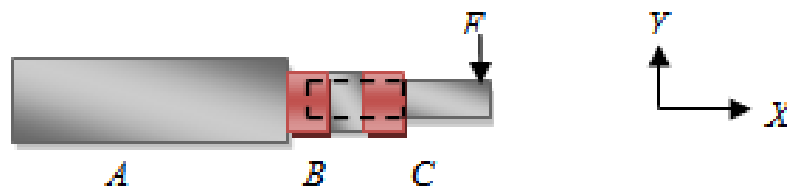


Figure 1: A Basic Schematic Illustration of Cutting Device Framework (A) Spindle, (B) Periodic Instrument Holder, and (C) Cutting End

As a rule, when a wave proliferating in a structure experiences an adjustment in the geometry and additionally the material properties, the wave is part into two segments; a spreading segment and a reflected segment. The reflected part collaborates with the proliferating part in a way that is controlled by the stage contrast between them

Investigations of the attributes of one-dimensional periodic structures have been widely revealed by Y. Altintas[9]. These structures are anything but difficult to dissect because of its geometrical straightforwardness. Yang [10] developed a systematic model adaptation methodology in order to continuously update the thermal-error model under varying manufacturing conditions. Process-periodic probing and adaptive system identification techniques are integrated to monitor and estimate machine-tool errors and recursively modify model coefficients as manufacturing process proceeds. Brown [11] presented a method is for determining the wave numbers, wave shapes and point reacceptances for an infinite, one-dimensional, non-uniform periodic structure with distributed periodic attachments or supports. The approach is based on a general theory of harmonic wave propagation in one-dimensional periodic systems. Periodic structures consist of an

arrangement of coupled identical substructures. When, due to unavoidable defects, discrepancies occur between the substructures, the periodicity is destroyed and the vibration localization phenomenon arises, consisting of a vibrational energy confinement in small regions of the structure, P. Bisegna [12]. Dong Li [13] presented a brief review of linear waves and dynamic behavior of both periodic and near-periodic structures. Al Ba'ba'a [14] achieved the stop band behavior via Bragg scattering in photonic media is most commonly evaluated using wave propagation models which predict gaps in the dispersion relations of the individual unit cells for a given frequency range.

The impacts of the excitation point, and additionally the flexible help attributes on the pass and stop qualities of the holder, are exhibited. The tool holder has a fastening side with a fastening projection and, facing away from said fastening side, an exterior, and during use of the tool, centrifugal forces are effective in the direction from the fastening side to the exterior. This invention ensures the free rotatability of a tool inserted in the tool holder even if overburden material enters the tool seat. For this purpose, the holding projection has an opening penetrating the inner wall of the tool seat and creating a spatial connection to the exterior, and the opening opens the tool seat towards the exterior.

The vast majority of tool condition monitoring systems use the cutting force as the predictor signal, Bhattacharyya [15]. Planewave propagation in infinite two-dimensional periodic lattices is investigated using Floquet-Bloch principles, Srikantha Phani [16]. Narisetti [17] investigated wave propagation in one-dimensional nonlinear periodic structures through a novel perturbation analysis and accompanying numerical simulations. Several chain unit cells are considered featuring a sequence of masses connected by linear and cubic springs. Brian R [18] described a method by which the dispersion relations for a two-dimensional structural component can be predicted from a finite element (FE) model. The structure is homogeneous in two dimensions but the properties might vary through the thickness. This wave/finite element (WFE) method involves post-processing the mass and stiffness matrices, found using conventional FE methods, of a segment of the structure. Y Yong [19] described a new method for the analysis of long and complicated structures which are composed of spatially periodic units or sections of spatially periodic units. The response of such a structure to external excitations is treated as a superposition of wave motions, with account taken of the effects of wave reflection due to change in the construction pattern along the structure and boundary conditions. Zhengyou Liu [20] extend the multiple-scattering theory for elastic waves by taking into account the full vector character. The formalism for both the band structure calculation and the reflection and transmission calculations for finite slabs is presented.

The vibration of cutting tool framework under specific conditions has for quite some time been perceived as a standout amongst the most critical elements influencing the execution of a machine device. Before, a few strategies for the distinguishing proof of processing vibrations have been proposed. The utilization of the steadiness graphs are considered. It has been demonstrated that the complex scientific estimations for processing elements dependent on a lot of slicing forces information are required to foresee the beginning of vibration by utilizing the dependability graphs. A reproduction to anticipate cutting powers and apparatus diversion amid end processing activity proposed, and to check the exactness of recreation results contrasted and those dependent on the hypothetical connections.

TRANSFER MATRIX ANALYSIS

The exchange matrix approach, mostly, depends on building up a connection between two finishes of the structure component. The genuine intensity of the exchange matrix approach comes, when the structure can be partitioned into an arrangement of substructures with an arrangement of components and hubs that are associated with another set on some

imaginary limit inside the structure. Utilizing the strategy for static buildup, the inner hubs/degrees of opportunity of the substructure can be dispensed with, along these lines decreasing the extent of the worldwide lattices of the structure.

When a set of equations for structural problems, can be manipulated to collect the forces and displacements of one end of the substructure on one side of the equation and relate them to those on the other end with a matrix relation, that matrix is called the transfer matrix of the structure. The transfer matrix of a substructure is then multiplied by that of the neighboring structure, in contrast with the superposition that is used in conventional numerical methods. Thus, the matrix system that describes the dynamics of the structure becomes significantly smaller in size. The transfer matrix method becomes of even more appealing features when identical substructures can be selected, thus, calculating the transfer matrix for one substructure is enough to describe all the dynamics of the whole structure. This particular feature is one that is inherent in all periodic structures by definition.

At the point when an arrangement of conditions for basic issues can be controlled to gather the powers and relocations of one end of the substructure on one side of the condition and relate them to those on the opposite end with a lattice connection, that framework is known as the transfer matrix of the structure. The transfer matrix of a substructure is then increased by that of the neighboring structure, interestingly with the superposition that is utilized in regular numerical techniques. Along these lines, the network framework that depicts the elements of the structure turns out to be essentially littler in size. The transfer matrix technique happens to considerably more engaging highlights when indistinguishable substructures can be chosen, along these lines, figuring the exchange framework for one substructure is sufficient to depict every one of the elements of the entire structure. This specific component is one that is natural in every single periodic structure by definition.

The examination of the periodic structures was drawn closer by various techniques; by far most of writing connected the exchange network approach. The got transfer matrix is portrayed by being shortsighted when gotten from a symmetric, traditionalist or non-preservationist, dynamic solidness framework. The essential property of a shortsighted lattice is that its eigenvalues show up in sets, one of which is the complementary of the other. This property of the transfer matrix has been taken a gander at as one that presents effortlessness for the examination; sadly, that equivalent property causes the numerical dangers in the investigation of structures with an extensive number of cells.

The spectral element matrix, often named the dynamic stiffness matrix, is known to provide the accurate dynamic characteristics of a structure because it is formed by exact shape functions, Usik Lee [21]. The proposed model was utilized for a structure that could be separated into substructures as strips whose hubs can be sorted out into two sets each lie on one side of the substructure. Because of the unpredictability of the coupling between adjoining cells in two-dimensional structures, the exchange network approach is not completely appropriate. In this examination, the connection between the outcomes acquired from the transfer matrix approach and those introduced by the spread surfaces will be concentrated to get a superior comprehension of the proliferation surfaces. Likewise, an endeavor to create engendering and lessening bends to portray the dynamic attributes of occasional plates will be presented. Static and dynamic distortions of machine tool holder assume a critical role in a machining procedure, which influencing the quality and profitability. Over the top prattle (self-energized vibration) may cause resistance infringement. Cutting force model can be utilized to anticipate and conquer these issues.

In this examination, summed up conditions are displayed which can be utilized for anticipating the static and dynamic properties of processing framework parts. Because of its wide use in industry, processing process is considered, be that as it may, similar techniques can be connected to other machining tasks also. Demonstrating of the processing procedure has been the subject of numerous investigations some of which are condensed. The focal point of these investigations has for the most part been on the displaying of cutting geometry, Bao [22]. In the present study, periodic elements are considered because the elements exhibit unique dynamic characteristics that make them act as mechanical filters for wave propagation, Nouh [23]. As a result, waves can propagate along the periodic elements only within specific frequency bands called the 'pass bands' and wave propagation is completely blocked within other frequency bands called the 'stop bands'. The ability of periodic structures to transmit waves from one location to another within the passbands, can be greatly reduced when the ideal periodicity is disrupted resulting in the well-known phenomenon of localization.

In the present investigation, periodic components are considered in light of the fact that these components show interesting unique attributes that make them go about as mechanical channels for wave engendering. Accordingly, waves can proliferate along the periodic components just inside explicit recurrence groups called the 'pass bands' and wave spread is totally hindered inside other recurrence groups called the 'stop bands'. The capacity of periodic structures to transmit waves starting with one area then onto the next inside the pass groups, can be extraordinarily lessened when the perfect periodicity is disturbed bringing about the notable marvel of limitation. Think about the vibration of a flexible tool holder, of length l and having an unvarying round cross-sectional territory in the xy -bearing typical to the z -pivot for vertical processing task (figure 2). The vibration of the holder can be displayed as a pole with one end at the base. The impacts of the shaft engine are represented by including their aggregate dormancy.

Usik [24] introduced an Fast Fourier transforms (FFT)-based spectral analysis method is for the dynamic analysis of the linear discrete dynamic system subjected to non-zero initial conditions. To evaluate the proposed FFT-based spectral analysis method, the forced vibration of a three degree-of-freedom (DOF) system is considered as an illustrative problem.

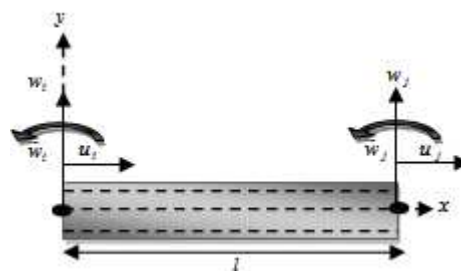


Figure 2: Straight Material Cutting Tool Holder

EQUATION OF MOTION AND BOUNDARY CONDITIONS

Since there are six nodal factors for the holder component, four for bending and two for the pivotal powers, a cubic polynomial function is assumed for $w(x)$, and the first order for $u(x)$. To consider the element which has three components at each end, w_i , w_i' and u_i at the top of the holder, w_j , w_j' and u_j at the bottom parallel to the surface of the machine table. For constant values of EI and EA equation (1) may be integrated to yield equations (2), where, C_i are constants of integration respect to x .

$$\begin{bmatrix} w_i \\ \dot{w}_i \\ w_j \\ \dot{w}_j \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & l & l^2 & l^3 \\ 0 & 1 & 2l & 3l^2 \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{bmatrix} = [A_w]\{c\} \quad \begin{bmatrix} u_i \\ u_j \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 1 & l \end{bmatrix} \begin{bmatrix} c_5 \\ c_6 \end{bmatrix} = [A_u]\{c\} \quad (1)$$

$$\begin{bmatrix} c_1 \\ c_2 \\ c_3 \\ c_4 \end{bmatrix} = \frac{1}{l^3} \begin{bmatrix} l^3 & 0 & 0 & 0 \\ 0 & l^3 & 0 & 0 \\ -3l & -2l^2 & 3l & -1l^2 \\ 2 & l & -2 & l \end{bmatrix} \begin{bmatrix} w_i \\ \dot{w}_i \\ w_j \\ \dot{w}_j \end{bmatrix} = [A_w^{-1}]\{w\}, \quad \begin{bmatrix} c_5 \\ c_6 \end{bmatrix} = \frac{1}{l} \begin{bmatrix} l & 0 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} u_i \\ u_j \end{bmatrix} = [A_u^{-1}]\{u_{ij}\} \quad (2)$$

Utilizing condition (2) to discover the shape functions $\{N\}$, where $\{N\} = [A^{-1}]\{w\}/N$. Substitution of $\{N\}$ values into the expressions of $w(x)$ and $u(x)$ yields the approximation of the mode shapes in the following equations.

$$w(x) = \left(1 - \frac{3x^2}{l^2} + \frac{2x^3}{l^3}\right)w_i + \left(x - \frac{2x^2}{l} + \frac{x^3}{l^2}\right)w_i' + \left(\frac{3x^2}{l^2} - \frac{2x^3}{l^3}\right)w_j + \left(\frac{x^3}{l^2} - \frac{x^2}{l}\right)w_j' \quad (3)$$

$$u(x) = \left(1 - \frac{x}{l}\right)u_i + \left(\frac{x}{l}\right)u_j \quad (4)$$

Potential and Kinetic Energies

Consider the energy associated with approximation given by the previous equations (3) and (4). The potential energy (PE) of the tool holder is non-dimensionalised by EI/l , will be expressible as:

$$PE_T = PE_w + PE_u = \frac{1}{2} \left[\int_0^l EI [w''(x)]^2 dx + \int_0^l EA [u'(x)]^2 dx \right] \quad (5)$$

Hence the vector $\{N'\} = \frac{\partial}{\partial x} \{N\}$, with entries $\{N_1'\}$ through $\{N_6'\}$, the main subordinates for conditions (3) and (4) will be:

$$w'(x) = \left(\frac{6x^2}{l^3} - \frac{6x}{l^2}\right)w_i + \left(1 - \frac{4x}{l} + \frac{3x^2}{l^2}\right)w_i' + \left(\frac{6x}{l^2} - \frac{6x^2}{l^3}\right)w_j + \left(\frac{x^3}{l^2} - \frac{x^2}{l}\right)w_j' \quad (6)$$

$$u'(x) = \left(\frac{-1}{l}\right)u_i + \left(\frac{1}{l}\right)u_j \quad (7)$$

$$w''(x) = \left(\frac{12x}{l^3} - \frac{6}{l^2}\right)w_i + \left(\frac{6x}{l^2} - \frac{4}{l}\right)w_i' + \left(\frac{6x}{l^2} - \frac{12x}{l^3}\right)w_j + \left(\frac{6x}{l^2} - \frac{2}{l}\right)w_j' \quad (8)$$

Substitution of $\{N_w''\}$ and $\{N_u'\}$ values into the expression of $w''(x)$ and $u'(x)$ yields the estimation of equation (9):

$$PE_T^e = \frac{EI}{2} \int_0^l \left[\left(\frac{12x}{l^3} - \frac{6}{l^2}\right)w_i + \left(\frac{6x}{l^2} - \frac{4}{l}\right)w_i' + \left(\frac{6x}{l^2} - \frac{12x}{l^3}\right)w_j + \left(\frac{6x}{l^2} - \frac{2}{l}\right)w_j' \right]^2 dx + \frac{EA}{2} \int_0^l \left[\left(\frac{-1}{l}u_i + \frac{1}{l}u_j\right) \right]^2 dx \quad (9)$$

The last articulation can be perceived, as relative to the result of the transpose of the vectors w and u . assuming the holder rigidity EI and EA are constant within the elements. For each element, the wu -stiffness matrix K is:

$$K_w^e = \frac{EI}{l^3} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{bmatrix} \quad \& \quad K_u^e = \frac{EA}{l} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \quad (10)$$

That is:

$$PE_T = \frac{1}{2} [(w^T K w) + (u^T K u)] \quad (11)$$

The kinetic energy of the element (KE) can be written in the elective after frame:

$$KE = KE_w + KE_u = \frac{m}{2} \left[\int_0^l [\dot{w}(x)]^2 dx + \int_0^l [\dot{u}(x)]^2 dx \right] \quad (12)$$

Considering the maximum kinetic energy at the end part of the holder $KE = KE_{max}$.

$$\dot{w}(x)_w = \omega w(x) \text{ Moreover } \dot{u}(x)_u = \omega u(x)$$

$$KE = \frac{m\omega^2}{2} \left[\int_0^l \left[\left(1 - \frac{3x^2}{l^2} + \frac{2x^3}{l^3} \right) w_i + \left(x - \frac{2x^2}{l} + \frac{x^3}{l^2} \right) w_i' + \left(\frac{3x^2}{l^2} - \frac{2x^3}{l^3} \right) w_j + \left(\frac{x^3}{l^2} - \frac{x^2}{l} \right) w_j' \right]^2 dx + \int_0^l \left[\left(1 - \frac{x}{l} \right) u_i + \left(\frac{x}{l} \right) u_j \right]^2 dx \right] \quad (13)$$

For direct frameworks that comply with Rayleigh's correspondence guideline, Phani and Adhikari (2008), related the matrices M and K as pursues:

$$K - \omega^2 M = 0 \quad (14)$$

Where, ω is the natural frequency of an element. An eigenvalue examination must be performed in planning a basic framework that will be exposed to elements powers. By substituting an eigenvalue λ_i into equation (14):

$$[K - \lambda_i M] w_i = 0 \quad \& \quad [K - \lambda_i M] u_i = 0 \quad (15)$$

Where, eigenvectors w_i and u_i correspond to deflection mode that gives the shape of the element. Therefore, analysis of eigenvalue equations gives important information on possible deflection modes experienced by the structure when it undergoes forces. In equation (15), since the mass matrix (M) is symmetric positive definite and stiffness matrix (K) are symmetric and either positive or positive semi-definite, the eigenvalues are all real and either positive or zero. The corresponding eigenvalue equations are having multiple eigenvalues. For an eigenvalue of multiplicity N , there are N vectors satisfying equation (16). The kinetic energy relative to the displacement will be:

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$$KE_T = \frac{1}{2} [(w^T M_w w) + (u^T M_u u)] \quad (16)$$

Where, M is the mass matrix for the system elements and defined by:

$$M_w^e = \frac{m}{420} \begin{bmatrix} 156 & 22l & 54 & -13l \\ 22l & 4l^2 & 13l & -3l^2 \\ 54 & 13l & 156 & -22l \\ -13l & -3l^2 & -22l & 4l^2 \end{bmatrix} \quad \& \quad M_u^e = \frac{m}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \quad (17)$$

Using equations (10) and (16) the dynamic equations becomes:

$$\frac{m}{420} \begin{bmatrix} 156 & 22l & 54 & -13l \\ 22l & 4l^2 & 13l & -3l^2 \\ 54 & 13l & 156 & -22l \\ -13l & -3l^2 & -22l & 4l^2 \end{bmatrix} \begin{bmatrix} w_i'' \\ w_i''' \\ w_j'' \\ w_j''' \end{bmatrix} + \frac{EI}{l^3} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{bmatrix} \begin{bmatrix} w_i \\ w_i \\ w_j \\ w_j \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} \quad (18)$$

$$\frac{m}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{bmatrix} u_i'' \\ u_j'' \end{bmatrix} + \frac{EA}{l} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} u_i \\ u_j \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (19)$$

Prediction of Cutting Tool Holder Deflection

Numerous analytical methods are available to predict the stability of milling processes. Most of these methods base on the assumption, that the dynamics of the machine tool are time invariant, Brecher [25]. For milling structure (Akeson, 2009), only the tool/holder deflection during the cut-in process of each tooth will be imprinted directly on the machine surface. Therefore, concentrate the distribution of F_w during the cut-in process of each tooth will be considered. Basically, the cutting tool/holder deflection expressed as two degree of freedom system with their structural parameters (Kivanc, 2003). The equation of motion in the w and u direction expressed as:

$$m_w \ddot{w} + k_w w = F_w \quad \text{and} \quad m_u \ddot{u} + k_u u = F_u \quad (20)$$

Where (w, u) , (\dot{w}, \dot{u}) and (\ddot{w}, \ddot{u}) are the cutter displacement, velocities and accelerations in the w and u directions, respectively. m_w and k_w are the structural parameters in the w directions and m_u , and k_u are the structural parameters in the u directions. It is that the transfer function between the cutting forces on the cutter assumed linear and has a single degree of freedom with mass, damping ratio and natural frequency. In this simplified spindle system with a single degree of freedom, the deflection in the w and u directions are expressed as follows:

$$\begin{aligned} \ddot{w}_{i+1} &= F - k w_i / m & \ddot{u}_{i+1} &= F - k u_i / m \\ \dot{w}_{i+1} &= \dot{w}_i + \ddot{w}_{i+1} & \dot{u}_{i+1} &= \dot{u}_i + \ddot{u}_{i+1} \\ w_{i+1} &= w_i + \dot{w}_{i+1} dt & \text{And} & & u_{i+1} &= u_i + \dot{u}_{i+1} \end{aligned} \quad (21)$$

Prediction Function of Periodic Elements

One component for the cutting instrument holder display gives mistaken outcomes if higher modes are energized, in this manner, more components must be utilized to show the whole structure. In the event that different components are utilized, conditions for all components must be gathered into a model of whole structure overall. For dynamic examination

of the holder and consolidating conditions (4) and (5), the redirection is introduced inside a holder component as:

$$wu_{ij} = \left(1 - \frac{3x^2}{l^2} + \frac{2x^3}{l^3}\right)w_i + \left(x - \frac{2x^2}{l} + \frac{x^3}{l^2}\right)w'_i + \left(1 - \frac{x}{l}\right)u_i + \left(\frac{3x^2}{l^2} - \frac{2x^3}{l^3}\right)w_j + \left(\frac{x^3}{l^2} - \frac{x^2}{l}\right)w'_j + \left(\frac{x}{l}\right)u_j \quad (22)$$

MODEL OF VIBRATING TOOL HOLDER

It is helpful to have details of movement, which make utilization of amounts identifying with the entire framework from which components are made up. The conditions of movement can be acquired from the former articulations for dynamic KE and potential PE energies utilizing the variety or Lagrangian approach. Consolidating conditions (10) and (13) for getting the aggregate energies. The conditions of movement for the vibratory framework can be given in the structure as:

$$\sum_{e=1}^n [M_e] \{\ddot{\delta}_e\} + \sum_{e=1}^n [K_e] \{\delta_e\} = \sum_{e=1}^n \{F_e\} \quad (23)$$

$$[M] \{\ddot{\delta}\} + [K] \{\delta\} = \{F\} \quad (24)$$

The estimation of vibration communicated as:

$$\{\delta\} = [\{F\} - [M] \{\ddot{\delta}\} / [K]] \quad (25)$$

Where, $\{\delta_e\} = \{w_1 \ w'_1 \ u_1 \ \dots \ w_N \ w'_N \ u_N\}$ is nodal deflection vector of the element, n denoting number of nodal points and $\{F_e\}$ is the vector of external forces.

Taking the cutting tool holder path as circular arc moves by the feed per tooth (chip load), c , in case of up milling tooth 1 in position ϕ_1 engages over the arc of cut, where $\phi_s \langle \theta_1 \langle \theta_E$ and $\phi_2 = \phi_1 + \pi/2$. The force acting on the holder can be added and reflected into the F_u and F_w components in the tool axis:

$$F_u = K_s A_d c [\sin(\phi) \cos(\phi) + 0.3 \sin^2(\phi)] \quad (26)$$

$$F_w = K_s A_d c [\sin^2(\phi) - 0.3 \sin(\phi) \cos(\phi)] \quad (27)$$

$$F = \sum_{n=1}^N F_u + \sum_{n=1}^N F_w \quad (28)$$

Both force components are periodic in 2ϕ , where $\phi = 2\pi(N/60)t$.

Longitudinal Vibration

Consider three components model of the longitudinal vibration and with one degree of freedom as appeared in figure 3. Since the three cells of the framework with two unique materials mixes (spring steel-elastic and spring steel copper) are unbending and pivoting in the meantime with one edge. Every component of the model has an active and potential energy. The basic may take different structures for the instrument holder. The flexural unbending nature EI of the component must be considered. From condition (20) three sets of matrices and their comparing with indistinguishable conditions and distinctive arrangements of obscure modular relocations u_i , can be gathered together by superimposing them

to yield condition (29) in the frame:

$$\frac{m}{18} \begin{bmatrix} 2 & 1 & 0 & 0 \\ 1 & 4 & 1 & 0 \\ 0 & 1 & 4 & 1 \\ 0 & 0 & 1 & 2 \end{bmatrix} \begin{bmatrix} u_1'' \\ u_2'' \\ u_3'' \\ u_4'' \end{bmatrix} + \frac{3EA}{l} \begin{bmatrix} 1 & -1 & 0 & 0 \\ -1 & 2 & -1 & 0 \\ 0 & -1 & 2 & -1 \\ 0 & 0 & -1 & 1 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} \quad (29)$$

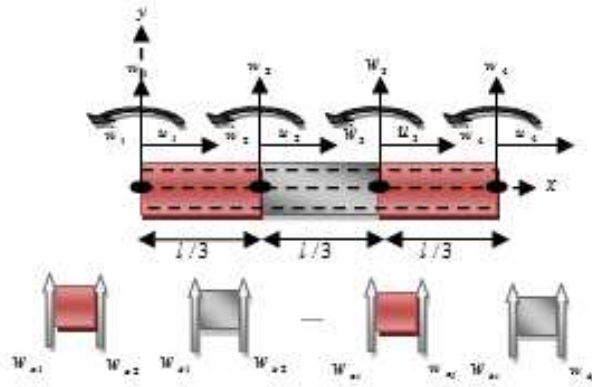


Figure 3: Periodic Cutting Tool Holder Model

Translation and Rotational Vibration

The mass and stiffness matrices for a braced holder at the shaft and free with the cutting instrument is appeared in figure 1. Using three elements and four nodes with $l=l/3$, the equations for the finite element at $(i=1, j=2)$, $(i=2, j=3)$, $(i=3, j=4)$, becomes:

$$\frac{m}{1260} \begin{bmatrix} 156 & \frac{22l}{3} & 54 & \frac{-13l}{3} \\ \frac{22l}{3} & \frac{4l^2}{9} & \frac{13l}{3} & \frac{-l^2}{3} \\ 54 & \frac{13l}{3} & 156 & \frac{-22l}{3} \\ \frac{-13l}{3} & \frac{-l^2}{3} & \frac{-22l}{3} & \frac{4l^2}{9} \end{bmatrix} \begin{bmatrix} \ddot{w}_i \\ \ddot{w}_i \\ \ddot{w}_j \\ \ddot{w}_j \end{bmatrix} + \frac{9EI}{l^3} \begin{bmatrix} 12 & \frac{2l}{9} & -12 & \frac{2l}{9} \\ \frac{2l}{9} & \frac{4l^2}{9} & -2l & \frac{2l^2}{9} \\ -12 & -2l & 12 & -2l \\ \frac{2l}{9} & \frac{2l^2}{9} & -2l & \frac{4l^2}{9} \end{bmatrix} \begin{bmatrix} w_i \\ w_i \\ w_j \\ w_j \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{bmatrix} \quad (30)$$

General structure of equations of motion is

$$M \begin{bmatrix} w_i'' \\ w_i''' \\ u_i'' \\ w_j'' \\ w_j''' \\ u_j'' \end{bmatrix} + K \begin{bmatrix} w_i \\ w_i \\ u_i \\ w_j \\ w_j \\ u_j \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{bmatrix}$$

General structure of conditions of movement for three periodic components:

$$M = \begin{bmatrix} M_c + M_a & M_b & 0 \\ M_b^T & M_c + M_a & M_b \\ 0 & M_b^T & M_c \end{bmatrix} \text{ and } K = \begin{bmatrix} K_c + K_a & K_b & 0 \\ K_b^T & K_c + K_a & K_b \\ 0 & K_b^T & K_c \end{bmatrix} \quad (31)$$

Where:

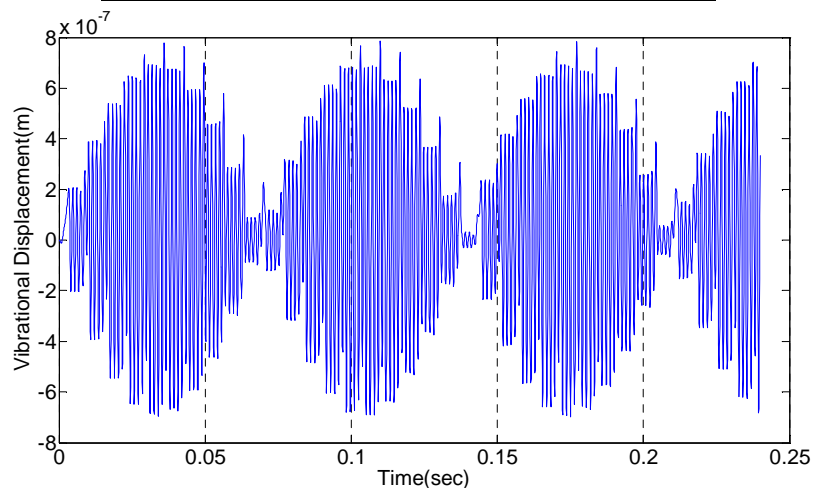
$$M_a = \begin{bmatrix} 156 & \frac{22l}{3} \\ \frac{22l}{3} & \frac{8l^2}{9} \end{bmatrix}, M_b = \begin{bmatrix} 54 & \frac{-13l}{3} \\ \frac{13l}{3} & \frac{-l^2}{9} \end{bmatrix}, M_c = \begin{bmatrix} 156 & \frac{-22l}{3} \\ \frac{-22l}{3} & \frac{8l^2}{9} \end{bmatrix}, K_a = \begin{bmatrix} 12 & \frac{2l}{9} \\ 2l & \frac{4l^2}{9} \end{bmatrix}, K_b = \begin{bmatrix} -12 & \frac{2l}{9} \\ -2l & \frac{2l^2}{9} \end{bmatrix}, \text{ and } K_c = \begin{bmatrix} 12 & \frac{-2l}{9} \\ -2l & \frac{4l^2}{9} \end{bmatrix}$$

SIMULATION RESULTS

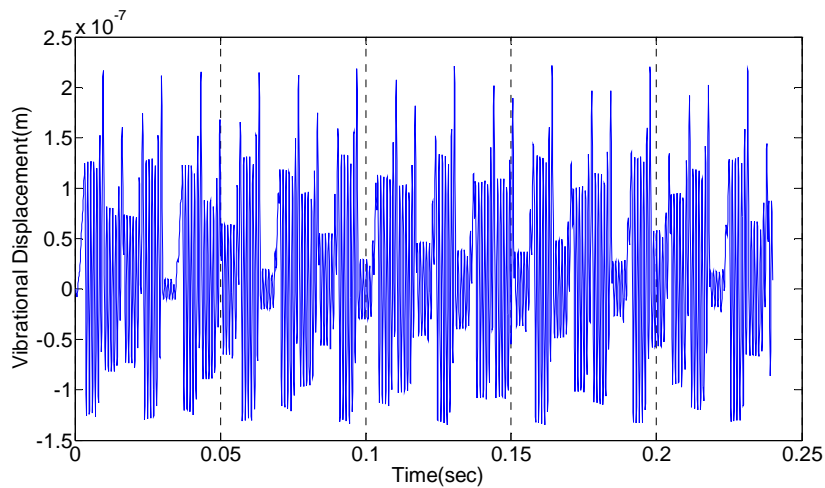
Know that a symphonious power produces consonant vibrations of a similar frequency, and the sufficiency of the vibrations relies upon the adequacy of the cutting power and on the proportion of the recurrence of the power over the normal recurrence of the framework. On the off chance that the two frequencies are equivalent, the instance of reverberation and greatest vibration adequacy will happen. In this paper, all simulations and representations depend on end process with helical smooth edges utilizing the proposed dynamic processing model with parameters recorded in table 1, in view of numerical hypothesis and strategy with the Eulerian approach (Jalili Saffar et al., 2008). The examination of the cutting power vibration and its impacts on constrained variety is plotted in the set graphs in figure 4, which contain time plots that delineate the strength enhancements from straight to periodic instrument holders, with four homogeneous teeth and different cutting commitment. The plots were gotten from the PC program written in MATLAB.

Table 1: Simulation Parameters

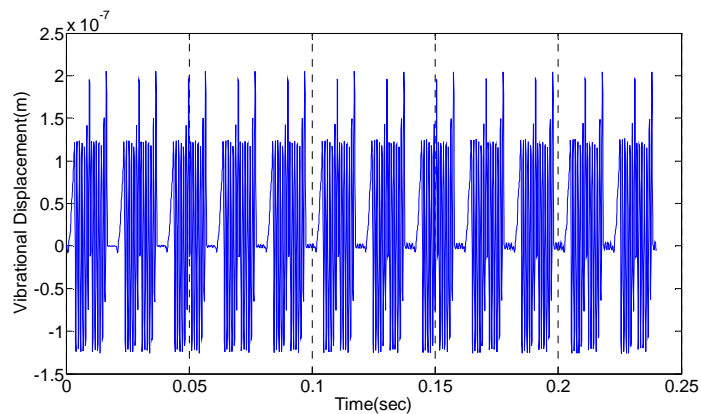
Specific Cutting Force, K_s	2100 N/mm
Nominal feed per tooth, f_T	2 mm
Cutting tool diameter, D	16 mm
Too holder diameter, D_h	24 mm
Number of cutting edges, Z	4
Axial depth of cut, A_d	5 mm
Radial depth of cut, R_d	3 mm
Spindle speed, N	3000 rpm
Number of samples, NS	1000
Modulus of Elasticity for Rubber, E	0.1 GPa
Modulus of Elasticity for Copper, E	117 GPa
Modulus of Elasticity for spring steel, E	210 GPa



(A)



(B)



(C)

Figure 4: Display Recreation Eventual Outcomes of Vibration Designs, (A) Periodic Spring Steel-Versatile, (B) Periodic Spring Steel-Copper and (C) Straight Spring Steel of Handling Instrument Holder

The program pursues the pivot of the cutter in 250 steps /upheaval, $d\theta = 360/240 = 1.44^\circ$, and it keeps running for 1000 stages, that is 4 revolutions. The tooth passing recurrence is NZ , where N signified the rotational speed and Z is the quantity of front lines of the processing cutter. The time between two back-to-back cuts (T) causes a stage distinction as $T = 1/NZ$. The feed per tooth (f_T) combined with a variable shaft speed (N) in a changing feed rate (f) which causes adjustment of the cutting powers F_u and F_w . Amid a few introductory tooth periods, vibrations begin to create and after that achieve the consistent state in which the vibration at the aggregate time (t) is resolved, where $t = 0.25$ sec, as appeared in the accompanying figure 4.

The resultant cutting power of all simulations under a similar cutting parameters of table 1, are in great assentation, as appeared in figure 5. Trusting that vibration will settle search for the cutting power in the last unrest, i.e. the last 250 stages, only to plot this part, as appeared in figure 6.

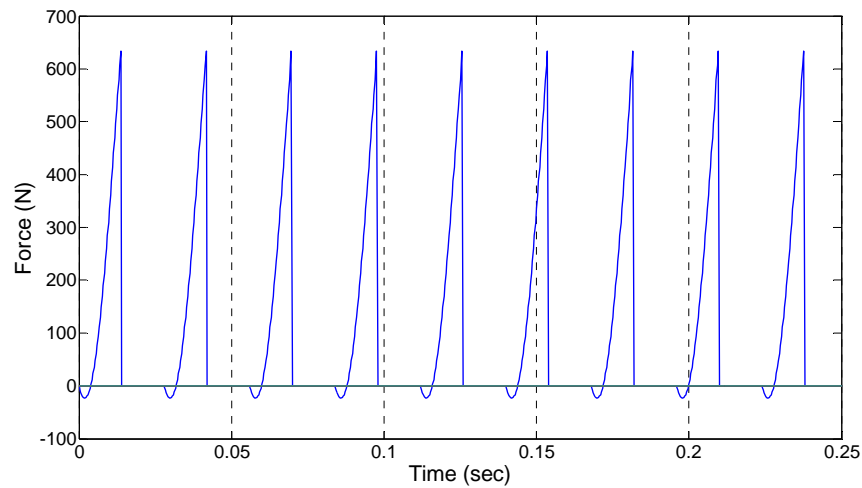


Figure 5: Milling Component Force on Cutter with 4 Straight Teeth and 0.25 Sec Cutting Time

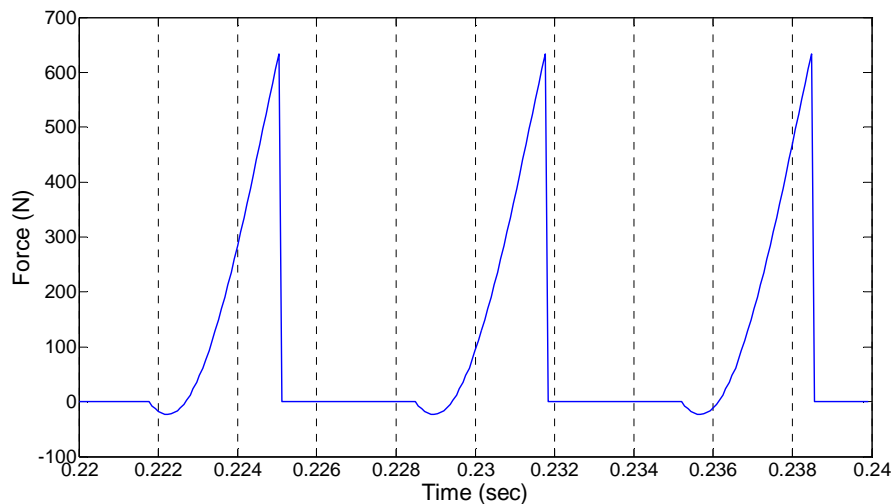


Figure 6: Force in Steady State for Last Revolution

CONCLUSIONS

Vibration is largely dodged either by solidifying the relative consistence between the cutting instrument framework and the workpiece, or by decreasing the pivotal and outspread profundities of cut. In this paper, another methodology for observing vibration amid the machining procedure by controlling the periodic materials of the machine apparatus holder is displayed. A computerized dynamic reenactment display was proposed to examine the impact of periodic cutting periodic holders and additionally basic parameters on the steadiness of processing vibrations. The model written in MatLab incorporates the commitment of the mass and firmness and its impact on the cutting power amplitudes. The paper displays another class of periodic machine device holder framework for detaching the vibration transmission from the slicing instrument holder to the machine device table trying to create a tranquil surface wrap up. A hypothetical model is created to portray the elements of wave engendering in an periodic apparatus holder. The model is inferred utilizing the hypothesis of limited components. The model of three periodic components, spring steel-elastic, and spring steel-copper and straight spring steel to process the vibration amplitudes and powers are introduced. The transfer matrix

detailing for every component is given. A correlation between those hypothetical methodologies with genuine estimations will be examined in the following examination.

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NOMENCLATURE

A: cross-sectional area, m^2 ; **A_a:** axial depth of cut, mm; **b:** chip width, mm; **C:** cutting force coefficient; **E:** Young's modulus, N/m²; **EI:** holder rigidity; **F:** total force magnitude, N; **F_u:** axial force, N; **F_w:** bending force, N; **f:** transverse force, N/m; **f_T:** feed per tooth, mm/tooth; **h:** undeformed chip thickness, mm; **I:** second moment of area, m⁴; **K:** coefficients of stiffness matrix; **KE:** kinetic energy, kgm²/s²; **l:** element length, m; **m:** equivalent mass per unit length, kg/m; **M:** coefficients moment of inertia, Nm; **N:** spindle speed, rev/s; **PE:** potential energy, kgm²/s²; **R_d:** radial depth of cut, **A:** cross-sectional area, m²; **A_a:** axial depth of cut, mm; **b:** chip width, mm; **C:** cutting force coefficient; **E:** Young's modulus, N/m²; **EI:** holder rigidity; **F:** total force magnitude, N; **F_u:** axial force, N; **F_w:** bending force, N; **f:** transverse force, N/m; **f_T:** feed per tooth, mm/tooth; **h:** undeformed chip thickness, mm; **I:** second moment of area, m⁴; **K:** coefficients of stiffness matrix; **KE:** kinetic energy, kgm²/s²; **l:** element length, m; **m:** equivalent mass per unit length, kg/m; **M:** coefficients moment of inertia, Nm; **N:** spindle speed, rev/s; **PE:** potential energy, kgm²/s²; **R_d:** radial depth of cut (mm); **t:**

time, s ; x : axial co-ordinate of the beam holder, m ; v : transverse displacement of the beam, m ; w : i^{th} natural frequency of the beam, rad/s ; Z : number of cutting edges; θ : cutting force angle; ρ : material mass density, kg/m^3 ; γ : torsional constant.